# Molecular Complexes between 1,2,4,5-Tetracyanobenzene and Some Aromatic Electron Donors 

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#### Abstract

The charge-transfer (CT) complexes of 1,2,4,5-tetracyanobenzene (TCNB) with substituted benzenes were studied by measuring electronic absorption spectra in solution and in crystal. Two CT absorption bands were observed for the TCNB complexes with mesitylene, hexamethylbenzene, dimethylaniline, and $\mathrm{N}, \mathrm{N}, \mathrm{N}^{\prime}, \mathrm{N}^{\prime}$ -tetramethyl-p-phenylenediamine. The separation between the two CT bands was found to be almost constant for these TCNB complexes. From this together with the other experimental and theoretical facts, it was concluded that the appearance of the two CT bands is due to the two closely located vacant orbitals of TCNB. From the measurement of the polarized absorption spectrum of the single crystal, the $31,200-\mathrm{cm}^{-1}$ band of TCNB was found to shift toward shorter wavelengths by $3600 \mathrm{~cm}^{-1}$ by the complex formation with mesitylene. This blue shift can be explained by the interaction of the locally (within TCNB) excited configuration with the CT configurations. A theoretical study was made by considering the configuration interaction among ground, CT, and locally excited configurations. It was theoretically demonstrated that the center of one of the benzene rings is located above the carbon atom or $\mathrm{C}-\mathrm{C}$ bond of the other benzene ring in the stable geometrical configuration.


Recently 1,2,4,5-tetracyanobenzene (TCNB) has been found to be one of typical electron acceptors and to form stable charge-transfer (abbreviated hereafter to CT) complexes ${ }^{2}$ with aromatic compounds. ${ }^{3-5}$ Zweig, et al. ${ }^{3}$ measured the absorption spectra of CT complexes of polycyclic aromatic hydrocarbons with TCNB and found a linear relationship between the transition energies corresponding to the CT absorption maxima and the highest occupied molecular orbitals of the donors. Bailey, et al., ${ }^{4}$ obtained the CT absorption maxima and the association constants for the TCNB complexes with pyrene and dimethylaniline in chloroform solution. Foster and Thomson ${ }^{5}$ determined the association constants and stabilization energies of $\mathrm{N}, \mathrm{N}, \mathrm{N}^{\prime}, \mathrm{N}^{\prime}$-tetramethyl- $p$-phenylenediamine (TMPD) complexes with polycyanobenzenes in chloroform and compared the results with those of similar polynitrobenzene complexes.
In the present paper, spectrophotometric studies have been made with CT complexes of various substituted benzenes with TCNB, special attention being paid to the appearance of two CT bands due to the closely located two vacant orbitals of the electron acceptor, and to the change in the absorption spectrum of TCNB itself. The polarized absorption spectrum of the single crystal was measured with the TCNBmesitylene complex, in order to obtain some knowledge concerning these two points. Theoretical consideration has also been carried out with some TCNB complexes by a method combining LCAO-MO with configuration interaction. The stable geometrical configuration and the observed spectrophotometric behaviors of the complexes were discussed on the basis of the theoretical results.
(1) Department of Chemistry, Faculty of Science, Nagoya University
(2) R. S. Mulliken, J. Phys. Chem., 56, 801 (1952); J. Am. Chem. Soc., 74, 811 (1952).
(3) A. Zweig, J. E. Lehnsen, W. G. Hodgson, and W. H. Jura, ibid., 85, 3937 (1963).
(4) A. S. Bailey, B. R. Henn, and J. M. Langton, Tetrahedron, 19, 161 (1963).
(5) R. Foster and T. J. Thomson, Trans. Faraday Soc., 59, 2287 (1963).

## Experimental Section

TCNB, prepared by Ishitani and Maruyama of our laboratory, ${ }^{6}$ was purified by recrystallization from ethanol. The melting point, $258^{\circ}$ in sealed tube (turning blue), agrees well with the value in the literature. ${ }^{7}$
Hexamethylbenzene (HMB) and TMPD were purified by recrystallization followed by sublimation. Benzene, toluene, $m$ xylene, and mesitylene were shaken with concentrated sulfuric acid, washed with water, dried over calcium chloride, and finally distilled. Aniline and $o-, m-$, and $p$-toluidine were purified by the vacuum distillation. Methylene chloride and tetrachloroethane, which were used as solvents for spectral measurements, were dried over calcium chloride and carefully distilled.
The absorption spectra of the solutions were measured with a Cary recording spectrophotometer Model 14 M . Solutions containing components were mixed separately with each other immediately before measurement. The concentrations of electron donors were taken to be in excess up to $20-100$ times that of TCNB. The temperature of the cell compartment of the spectrophotometer was kept constant by a heater with a thermister. A microspectrophotometer constructed by Tanaka ${ }^{8}$ was used for the polarized ultraviolet absorption measurement of single crystals of the TCNBmesitylene complex which was obtained from the TCNB solution of mesitylene. The crystal showed a rectangular face, and the axis of elongation coincides with an extinction position. This means that the elongated axis or the axis perpendicular to it is the crystalline $b$ axis in a monoclinic crystal or one of the axes in an orthorhombic crystal.

## Results and Discussion

The Absorption Spectra of the Solutions Containing TCNB and Some Electron Donors. New absorption bands due neither to electron donor nor electron acceptor appear in the near-ultraviolet and visible regions for the mixed solution containing TCNB and such aromatic molecules as HMB, N,N-dimethylaniline, and TMPD, which are shown in Figure 1. For example, the TCNB-TMPD system ${ }^{9}$ shows two new bands at 13,900 and $20,300 \mathrm{~cm}^{-1}$. These bands may be regarded as the CT bands characteristic of the interaction between electron donor and acceptor.
(6) The authors wish to thank Mr. A. Ishitani and Miss Y. Maruyama for offering TCNB.
(7) E. A. Lawton and D. D. McRichie, J. Org. Chem., 24, 26 (1959).
(8) J. Tanaka, Bull. Chem. Soc. Japan, 36, 833 (1963).
(9) TCNB itself has the absorption band at $31,600 \mathrm{~cm}^{-1}$ in tetrachloroethane and TMPD itself at $31,550 \mathrm{~cm}^{-1}$.


Figure 1. Spectra of the TCNB complexes with (1) TMPD in $\mathrm{C}_{2} \mathrm{H}_{2} \mathrm{Cl}_{4}$; (2) dimethylaniline in $\mathrm{C}_{2} \mathrm{H}_{2} \mathrm{Cl}_{4}$; (3) hexamethylbenzene in $\mathrm{CH}_{2} \mathrm{Cl}_{2}$.

As for the benzene-TCNB system, a new absorption maximum is difficult to find by the usual procedure of dissolving the component molecules in tetrachloroethane. This is because the new band overlaps with that pertinent to TCNB itself (see Figure 2) and because this complex has a small equilibrium constant. Therefore, we dissolved TCNB in pure benzene and found that the new absorption maximum appears at 32,500 $\mathrm{cm}^{-1}$ (see Figure 2). The absorption bands were also measured by dissolving TCNB in pure donors like toluene, $m$-xylene, and mesitylene. The results are shown in Figure 2.

The peak wavelengths observed by the present authors are summarized in Table I. From the results given in this table, it may be inferred that the maximum wave-

Table I. CT Bands of the $\pi$-Complexes of TCNB with Some Aromatic Donors

${ }^{a}$ Absorption maximum in the pure liquid of donor. ${ }^{b}$ Foster and Thomson ${ }^{5}$ reported $\nu_{1} 23,400,24,700,18,600$, and 13,600 $\mathrm{cm}^{-1}$ for the HMB, naphthalene, $\mathrm{N}, \mathrm{N}$-dimethylaniline, and TMPD complexes, respectively. Furthermore, they found the second CT band at $19,600 \mathrm{~cm}^{-1}$ for the TMPD-TCNB complex. Their measurements were made with the chloroform solutions. ${ }^{\circ}$ Taken from ref 4. d Absorption maxima measured with the single crystal.
lengths of the new bands shift toward longer wavelengths with the lower ionization potentials of the donors. In order to check this point more quantita-


Figure 2. Spectra of TCNB in various solvents: (1) in $\mathrm{C}_{2} \mathrm{H}_{2} \mathrm{Cl}_{4}$, (2) in benzene, (3) in toluene, (4) in $m$-xylene, (5) in mesitylene.


Figure 3. Plot of $\nu_{\max }$ for the TCNB complexes with various donors against those for the corresponding 1,3,5-trinitrobenzene complexes.
tively, we plotted the maximum frequencies of the new bands observed with TCNB complexes against those observed with $1,3,5$-trinitrobenzene (TNB) complexes. As is shown in Figure 3, the points fit well to two straight lines. One of them (line I) corresponds to the first CT bands and the other (line II) to the second CT bands.
The Polarized Absorption Spectrum of the TCNBMesitylene Crystal. The polarized absorption spectrum of the TCNB-mesitylene single crystal was measured in the direction parallel ( $\mid$ ) and perpendicular $(\perp)$ to the elongated axis of the crystal. The results are shown in Figure 4. As is easily seen in Figure 4, three absorption peaks were found at $28,300,33,700$, and $34,800 \mathrm{~cm}^{-1}$, and a strong absorption tail was observed in the higher-frequency region than 37,000 $\mathrm{cm}^{-1}$.
The first absorption peak ${ }^{10}$ observed with the single crystal undoubtedly corresponds to the first CT band which appears at $28,400 \mathrm{~cm}^{-1}$ in solution. The absorption intensity is stronger in the direction parallel to the elongated axis of the crystal than in the perpendicular direction. Since the CT band may be
(10) Figure 4 shows that this band seems to split by $700 \mathrm{~cm}^{-1}$ in the single crystal. The splitting cannot be observed with the solution.


Figure 4. Polarized absorption spectrum observed with the single crystal of the TCNB-mesitylene complex: ——, with the polarized light parallel to the elongation axis; ----..---, with the polarized light perpendicular to the elongation axis. The vertical lines $E_{1}, \ldots, E_{5}$ show the calculated transition energies and oscillator strengths.
polarized perpendicular to the benzene rings of the electron donor and acceptor, the observed dichroism of the first band shows that the donor and the acceptor are stacked with each other in the single crystal in such a way that the benzene ring planes are almost perpendicular to the elongated axis. ${ }^{11}$

The second and third absorption bands observed with the single crystal may conceivably correspond to the absorption tail observed with the solution in the higher-frequency region than $31,000 \mathrm{~cm}^{-1}$. The intensity of the second band is stronger in the direction parallel to the elongated axis of the crystal. Therefore, the band may be regarded as the second CT band. The separation between the first and the second bands is $5400 \mathrm{~cm}^{-1}$. This agrees considerably well with the separation between the two CT bands observed with the TCNB-HMB and TCNB-TMPD complexes in solution. This point will be discussed later.

The dichroic intensity ratio of the first CT band is somewhat different from that of the second CT band. That is to say, the observed dichroic ratio ( $D_{i \mid}: D_{\perp}$ ) is 2.78 and 1.15 for the first and second bands, respectively. This will be explained by the mixing of the CT configurations with the locally excited configurations. In fact, the dichroic ratio ( $M_{2}{ }^{2} /\left(M_{x}{ }^{2}+M_{y}{ }^{2}\right)$ ) calculated by considering the configuration interaction between the CT and locally excited configurations is 3.47 and 0.87 for the first and second bands, respectively.

The dichroic intensity relation of the third band at $34,800 \mathrm{~cm}^{-1}$ is the reverse of those of the above-mentioned two bands. This means that the band is polarized in the direction parallel to the ring plane,

[^0]

Figure 5. Polarized absorption spectrum of the TCNB single crystal: -- with the polarized light perpendicular to the elongation axis; ..... with the polarized light parallel to the elongation axis.
and therefore it may be regarded as a shifted band of the $31,200-\mathrm{cm}^{-1}$ band of TCNB. The spectrum of the single crystal of TCNB itself was measured for the purpose of obtaining a quantitative knowledge about the shift. The result given in Figure 5 shows that the absorption maximum of the single crystal appears at $31,200 \mathrm{~cm}^{-1}$, i.e., at the same position as that of the solution spectrum. Thus it is concluded that the band of TCNB shifts toward shorter wavelengths by the complex formation and the magnitude of the shift amounts to $3600 \mathrm{~cm}^{-1}$.

Appearance of the Two CT Absorption Bands. The absorption peaks of the two CT bands were clearly observed for the TCNB complexes with HMB, TMPD, and $\mathrm{N}, \mathrm{N}$-dimethylaniline (Figure 1). ${ }^{12}$ A shoulder due to the second CT transition was observed for the TCNB complexes with naphthalene, aniline, and $o-$, $m$-, and $p$-toluidines.

Multi-CT absorption bands have been found for the CT complexes containing strong electron acceptors such as chloranil, bromanil, and tetracyanoethylene (TCNE). ${ }^{13}$ The origin of the multi-CT bands hitherto observed was explained by the existence of two closely located occupied orbitals of the donor. ${ }^{13}$ In this case, the frequency difference ( $\Delta \nu_{\mathrm{CT}}$ ) between the two CT bands is approximately equal to the energy difference in the first and the second highest occupied orbitals of the donor. The $\Delta \nu_{\mathrm{CT}}$ values for the chloranil and TCNE complexes, which are given in the third and fourth columns of Table II, respectively, are nearly coincident with each other for the respective donors.

The $\Delta \nu_{\text {CT }}$ values for the TCNB complexes are summarized in the second column of Table II. The $\Delta \nu_{\mathrm{CT}}$ value for the TCNE-TMPD complex is greater than twice that of the TCNB-TMPD complex. This seems to mean that the appearance of the two CT bands for the TCNB complexes cannot be explained
(12) There is a possibility of observing CT bands more than two for the TCNB complex with TMPD or N,N-dimethylaniline, because the CT band splits by the two ionization potentials of the electron donor as well as by the two electron affinities of TCNB. The third CT band, however, could not actually be observed. This is because it may appear in the shorter wavelength side than the measurable region.
(13) G. Briegleb, "Elektronen-Donator-Acceptor-Komplexe," Springer-Verlag, Berlin, 1961.

Table II. The Frequency Difference ( $\Delta \nu_{C_{T}}$ ) between the Two CT Bands ( $\mathrm{cm}^{-1}$ )

|  |  |  |  |
| :--- | :--- | :---: | :---: |
|  | Donors | TCNB | TCNE $^{a}$ | | Chlor- |
| :---: |
| anil |

${ }^{a}$ Taken from ref 13. ${ }^{b}$ Obtained with the single crystal.
by the same mechanism as that for the TCNE or chloranil complexes. In this connection, it is to be noticed that the $\Delta \nu_{\mathrm{CT}}$ values for the TCNB complexes are almost constant ( $6200 \mathrm{~cm}^{-1}$ ), independent from the electron donors. Furthermore, the observed $\Delta \nu_{\mathrm{CT}}$ values are very close to the calculated energy difference ( $\Delta \nu$ ) between the lowest and the second lowest vacant orbitals of TCNB; i.e., the observed $\Delta \nu_{\mathrm{CT}}$ and the theoretical $\Delta \nu$ are 6200 and $5500 \mathrm{~cm}^{-1}$, respectively. The details of the calculation of the $\Delta \nu$ value will be described in the theoretical part of the present paper. The above-mentioned facts lead to the conclusion that the two closely located vacant orbitals of the electron acceptor play an important role in the appearance of the two CT bands for the case of the TCNB complexes.

The above conclusion is also supported by the fact that the TCNB complex with hexamethylbenzene as an electron donor, which has the same symmetry as benzene and therefore whose two highest occupied orbitals are doubly degenerated, shows the two CT bands (see Figure 1). The present experimental result may be the first finding of the two CT bands due to the two closely located vacant orbitals of electron acceptor.

Molar Extinction Coefficients of TCNB Complexes. On the assumption that all the TCNB molecules form complexes when they are dissolved in pure liquids of benzene, ${ }^{14}$ toluene, and mesitylene, the peak molar extinction coefficients of the CT bands ( $\epsilon_{\max }$ ) can be obtained as is shown in Table III. In this table, the

Table III. Molar Extinction Coefficients $\left(\epsilon_{\max }\right)$ of the CT Bands

| Donors | $\overline{\mathrm{TCNB}}^{\epsilon_{\max }}$ | TNB ${ }^{\text {b }}$ | $I_{\mathrm{p}},{ }^{c} \mathrm{eV}$ |
| :---: | :---: | :---: | :---: |
| Benzene | 2200 ${ }^{\text {a }}$ | 5855 | 9.245 |
|  | 3300 (in tetrachloroethane) |  |  |
| Toluene | $1800^{\text {a }}$ | 3540 | 8.82 |
| Mesitylene | $1400{ }^{\text {a }}$ |  | 8.40 |
| HMB | 1000 (in methylene chloride) | 2150 | 7.85 |

${ }^{a}$ The values in pure liquid of donors. ${ }^{b}$ The values in $\mathrm{CCl}_{4}$ : G Briegleb and J. Czekalla, Z. Elektrochem., 59, 184 (1955). ${ }^{c}$ Ionization potentials of the donors obtained from ref 18.

[^1]

Figure 6a. Calculated transition moments: $\quad M_{z j^{2}} \equiv\left(\Psi \mathbf{G}|Z| \Psi_{j}\right)^{2}$. $(1+2)=M_{21}{ }^{2}+M_{22}{ }^{2} ;(3+4)=M_{23^{2}}{ }^{2}+M_{24}{ }^{2}$.


Figure 6b. Calculated transition moments: $\quad M_{j}{ }^{2} \equiv\left(\Psi \mathrm{G}|r| \Psi_{j}\right)^{2}$. $(1+2)=M_{1}{ }^{2}+M_{2}{ }^{2} ;(3+4)=M_{3}{ }^{2}+M_{4}{ }^{2}$.
$\epsilon_{\max }$ values for the TNB complexes were also shown for the purpose of comparison. The $\epsilon_{\text {max }}$ value for the HMB-TCNB complex was evaluated by the aid of the equilibrium constant, the determination of which will be described in a later part of the present paper. The $\epsilon_{\max }$ value of the benzene-TCNB complex was also evaluated by the aid of the equilibrium constant as well as from the data of the TCNB in pure benzene. As is seen in Table III, a discrepancy exists between these two $\epsilon_{\max }$ values. This might be due partly to the difference in the dielectric constants of the solvents.
Table III shows that the CT band intensity increases with the decreasing donor ability, i.e., with the higher ionization potentials ( $I_{\mathrm{p}}$ ) of the donors. This may be explained by considering the configuration interaction between the CT configuration and locally (within TCNB) excited configuration. The present theoretical consideration, the details of which will be described in the Theoretical Part, shows that the mixing of the locally


Figure 7. Calculated transition energies in $R_{x}=1.20, R_{y}=$ 0.695 , and $R_{z}=3.4 \AA$. The open circles represent the calculated transition energies for the TCNB complexes with various $I-A$ values. The straight lines 1 and 2 show the energies of the locally (within TCNB) excited configuration $\Phi_{\mathrm{L}!}$ and $\Phi_{L \varepsilon}$, respectively. The straight lines 3 and 4 show the energies of the CT configurations $\Phi_{\mathrm{CT}_{1}}$ and $\Phi_{\mathrm{CT}_{4}}$, and $\Phi_{\mathrm{Cr}_{2}}$ and $\Phi_{\mathrm{CT}_{4}}$, respectively. $a, b, c$, and d are respectively the $I-A$ values corresponding to the HMB , mesitylene, toluene, and benzene complexes.
(within the acceptor) excited configuration with the $C T$ configurations is important for the CT band intensity; that is to say, the second term of eq 13 in the Theoretical Part shows that the CT absorption bands borrow the intensity from the acceptor excitation. This effect becomes more important when the energy of the CT configuration $\left(E\left(\mathrm{D}^{+} \mathrm{A}^{-}\right)\right.$) becomes higher and is closer to the energy of the locally excited configuration ( $E\left(\mathrm{DA}^{*}\right)$ ). The theoretical relation between the donor ability and the CT band intensity is shown in Figure 6 b . The abscissa of this figure, $I-A$, represents the difference between the ionization potentials of the donors and the electron affinity of the acceptor, TCNB. Curve $(1+2)$ in the figure, which represents the square of the transition moment for the first CT band, rises steeply with the increment of the $I-A$ values. Thus it may be possible to explain the fact that the CT band intensity increases with the higher ionization potentials of the donors.

Orgel and Mulliken ${ }^{15}$ discussed the intensity borrowing for the contact-pair mechanism for TNB complexes. The present experimental and theoretical results indicate that the similar phenomenon is also important for the stable CT complex.

Murrell ${ }^{16}$ discussed the intensity borrowing of the CT bands from donor excitations. In the present study, we have shown that the intensity donation also occurs from acceptor excitations to the CT bands.

Comparison of the Observed and Theoretical Transition Energies. The transition energies of TCNB complexes were calculated by the present authors by considering the configurational interaction among the ground, two locally (within TCNB) excited, and four CT configurations. The transition energies evaluated

[^2]for various $I-A$ values are given in Figure 7. For the donor molecules with small ionization potentials, the two CT absorption bands corresponding to the transitions to the lowest four excited states ${ }^{17}$ are expected to appear in the visible and near-ultraviolet regions. The actual appearance of the two CT bands has already been discussed in the present paper. The transition energies of the HMB-TCNB complex were calculated to be 2.93 and 3.51 ev on the assumption of $I-A$ $=6.25 \mathrm{ev}$. These values may be said to agree well with the observed values, 2.91 and 3.65 ev . Since the ionization potential of HMB is $7.85 \mathrm{ev},{ }^{18}$ the electron affinity of TCNB was determined to be 1.60 ev. ${ }^{19}$

Using the above value of electron affinity and the ionization potential of mesitylene, $I=8.40 \mathrm{ev},{ }^{18}$ the transition energies for the mesitylene-TCNB complex were calculated to be 3.44, 4.06, and $4.25 \mathrm{ev} .{ }^{20}$ These values are well coincident with the transition energies observed with the crystal, $3.51,4.18$, and 4.34 ev . In Figure 4 , which shows the polarized absorption spectrum of the mesitylene-TCNB complex crystal, the calculated transition energies are shown by vertical lines, $E_{1}, E_{2}, E_{3}, E_{4}$, and $E_{5} . E_{1}, E_{2}, E_{3}, E_{4}$, and $E_{3}$ mean respectively the transition energies for the $W_{0} \rightarrow W_{1}, W_{0} \rightarrow W_{2}, W_{0} \rightarrow W_{3}, W_{0} \rightarrow W_{4}$, and $W_{0} \rightarrow$ $W_{5}$ transitions given in Table VIII in the Theoretical Part. The first CT band corresponds to $E_{1}$ and $E_{2}$ transitions and the second CT band to the $E_{3}$ and $E_{4}$ transitions. The doublet characters of these transitions are caused by the doubly degenerate highest occupied orbitals of mesitylene. The broadness of the first and the second CT bands may be due to the above-mentioned doublet characters. A similar calculation was also made for benzene and toluene complexes. The calculated transition energies to the lowest excited states are 3.70 and 3.55 ev for the benzene and toluene complexes, respectively, while the observed values are 4.03 and 3.94 ev .

Equilibrium Constants. The equilibrium constant ( $K$ ) for the benzene-TCNB complex was determined in tetrachloroethane. The concentration of benzene was varied between 0 and 5.05 mole/l., while that of TCNB was taken nearly constant (order of $10^{-4}$ mole/l.). TCNB has an absorption tail at $30,300 \mathrm{~cm}^{-1}$ where the equilibrium constant for the benzeneTCNB complex was measured. Therefore, we determined the equilibrium constant by the following equation which was derived by considering the change in the absorption spectrum of TCNB due to the complex formation at the wavelength under investigation ${ }^{21}$

$$
\begin{equation*}
\epsilon^{*}-\epsilon_{\mathrm{A}}=\frac{\epsilon-\epsilon_{\mathrm{A}}}{1+\frac{1}{K C_{\mathrm{D}}}} \tag{1}
\end{equation*}
$$

(17) These four excited states are separated into two sets, each of which consists of two closely located states.
(18) F. I. Vilesov, Soviet Phys. Usp., 6, 888 (1964); Usp. Fiz. Nauk, 81, 669 (1963).
(19) The electron affinity of TCNB is estimated as 0.4 ev by Briegleb. The value of the Coulomb attraction energy which was used in his review on the electron affinity (G. Briegleb, Angew. Chem. Intern. Ed. Engl., 3, 617 (1964)) is 4.3 ev . This value of the Coulomb attraction energy is apparently overestimated. Our Coulomb attraction energies are 3.30 $\sim 3.45$ ev as will be described later. Therefore his value of the electron affinity is smaller than ours by about 1 ev .
(20) The average values for the corresponding two transitions are 3.44 and 4.06 ev , respectively.
where $\epsilon^{*}\left(\equiv D / C_{\mathrm{A}} I\right)$ is the apparent molar extinction coefficient of the solution, $\epsilon_{\mathrm{A}}$ and $\epsilon$ are the molar extinction coefficients of TCNB and the complex, respectively, $C_{\mathrm{A}}$ and $C_{\mathrm{D}}$ are respectively the concentrations of TCNB and benzene, $D$ is the absorbance, and $l$ is the light path length of a cell in cm . The equilibrium constants were determined at $30,300 \mathrm{~cm}^{-1}$ using the least-squares method, and the results are given in Table IV.

Table IV. Equilibrium Constants ( $K$ ), Molar Extinction Coefficients ( $\epsilon$ ) of the CT Band, Heat of Formation ( $\Delta H$ ), and Entropy Change ( $\Delta S$ ) for the Benzene-TCNB Complex in $\mathrm{C}_{2} \mathrm{H}_{2} \mathrm{Cl}_{4}$

| Temp, <br> ${ }^{\circ} \mathrm{C}$ | $K$, <br> $1 . / \mathrm{mole}$ | $\boldsymbol{\epsilon}_{30,300 \mathrm{~cm}^{-1}}$ |
| :---: | :---: | :---: |
| 24 | 0.25 | 2400 |
| 28 | 0.23 | 2500 |
| 32 | 0.22 | 2500 |
|  | $\Delta H=-0.4 \mathrm{kcal} / \mathrm{mole}$ |  |
|  | $\Delta S=-0.9 \mathrm{eu}$ |  |

The equilibrium constant of the hexamethylbenzeneTCNB complex in methylene chloride was determined to be $3.0 \mathrm{l} . / \mathrm{mole}$ at $22.4^{\circ}$ by applying the BenesiHildebrand ${ }^{22}$ method to the absorbance values observed at $23,500 \mathrm{~cm}^{-1}$.

Heats of Formation. The heat of formation of the benzene-TCNB system was determined to be about 0.4 $\mathrm{kcal} /$ mole from the temperature dependence of the equilibrium constants given in Table IV. On the other hand, the heat of formation of the TMPD-TCNB complex was measured by Foster and Thomson to be about $3.6 \mathrm{kcal} / \mathrm{mole}^{5}$ in chloroform solution.

The stabilization energy due to the CT interaction was calculated by the method described in the Theoretical Part to be $0.03 \mathrm{ev}(0.7 \mathrm{kcal} / \mathrm{mole})$ and $0.06 \mathrm{ev}(1.4 \mathrm{kcal} /$ mole) for the ( $I-A$ ) values of 8.65 and 5.85 ev , respectively. The stabilization energy largely depends on the relative orientation of the electron donor and acceptor. The above values were evaluated for the stable conformation (model 2 described in the Theoretical Part).

The observed heat of formation of the complex may conceivably include the contributions of the electrostatic and dispersion forces in addition to the abovementioned CT force. Furthermore, as is well known, it is difficult to get an accurate result in the calculation of the stabilization energy, especially in this case, because of the ambiguity in the geometry of the complex. Therefore, the quantitative comparison of the calculated CT energy with the observed heat of formation might be rather meaningless. However, theoretical evalua-
(21) From the relation among $D, \epsilon_{\mathrm{A}}, \epsilon, C_{\mathrm{A}}$, and the concentration of the complex $C_{A D}$, we can obtain the following equation

$$
\begin{equation*}
D / C_{\mathrm{A}} l \equiv \epsilon^{*}=(1-x) \epsilon_{\mathrm{A}}+x \epsilon \tag{2}
\end{equation*}
$$

where $x=C_{A D} / C_{A}$. By solving this equation

$$
\begin{equation*}
x=\left(\epsilon^{*}-\epsilon_{A}\right) /\left(\epsilon-\epsilon_{A}\right) \tag{3}
\end{equation*}
$$

In the case of $C_{A} \ll C_{D}$, the equilibrium constant $K$ is related with $x$ as

$$
\begin{equation*}
K=C_{\mathrm{AD}} /\left[\left(C_{\mathrm{A}}-C_{\mathrm{AD}}\right) C_{\mathrm{C}}\right]=x /(1-x) C_{\mathrm{D}} \tag{4}
\end{equation*}
$$

By inserting eq 3 into eq 4 , eq 1 is easily derived.
(22) H. A. Benesi and J. H. Hildebrand, J. Am. Chem. Soc., 71, 2703 (1949).


Figure 8. Molecular orbitals and their symmetries of TCNB and benzene.
tion of the CT force may be helpful for the consideration of the stable conformation of the complex, because the CT force may conceivably be sensitive to the relative orientation of the donor to the acceptor, compared with the other forces.

## Theoretical Part

The electronic structures of the $\pi-\pi$ type CT complexes between benzene or methyl-substituted benzenes as electron donors and TCNB as electron acceptor were calculated by considering the configuration interaction among the ground (or no bond), CT (or dative), and locally excited configurations.

The electronic structures of the CT complexes are usually described in terms of the "no bond structure" and the "dative structure" as follows.

$$
\begin{gathered}
\Psi_{\mathrm{G}}=a \Phi(\mathrm{DA})+b \Phi\left(\mathrm{D}^{+} \mathrm{A}^{-}\right) \\
\Psi_{\mathrm{E}}=b^{\prime} \Phi(\mathrm{DA})-a^{\prime} \Phi\left(\mathrm{D}^{+} \mathrm{A}^{-}\right)
\end{gathered}
$$

Mulliken, ${ }^{23}$ however, pointed out that the other configurations like the back CT configurations and locally excited configurations are important in some cases. Hosoya and Nagakura ${ }^{24}$ demonstrated that the lowest excited state of the complex between cyclohexene and the silver cation has the mixing character of CT and donor Rydberg excitation. Furthermore, our experimental results suggest that the mixing of locally excited configurations $\Phi\left(\mathrm{D}^{*} \mathrm{~A}\right)$ and $\Phi\left(\mathrm{DA}^{*}\right)$ is important, in particular, in excited states. Here $\Phi\left(\mathrm{D}^{*} \mathrm{~A}\right)$ and $\Phi\left(\mathrm{DA}^{*}\right)$ mean the excited configurations caused by the excitation within the donor and the acceptor, respectively. In the present calculation we took into account the interaction among the ground, locally excited, and CT configurations.

First of all, the molecular orbitals of TCNB, $\phi_{1}$, $\phi_{2}, \ldots, \phi_{14}$ were calculated with the results given in Table V and Figure 8. The details of MO calculations are given in the Appendix. As for the molecular orbitals of methyl-substituted benzenes, we have used those of benzene which are represented by $\theta_{1}, \theta_{2}, \ldots$, $\theta_{6}$. This is because the small change in the shape of the MO may not give significant influence to the present results.

The wave function of the ground configuration is given by a Slater determinant

$$
\begin{equation*}
\Phi_{\mathrm{G}}=\left|\theta_{1} \vec{\theta}_{1} \cdots \theta_{3} \bar{\theta}_{3} \phi_{1} \bar{\phi}_{1} \cdots \phi_{7} \bar{\phi}_{7}\right| \tag{5}
\end{equation*}
$$

[^3]|  | Orbital energy, ev |  | W | n |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | $-16.594$ | $0.3108 \varphi_{0}{ }^{88}$ | $+0.3583 \varphi_{1}{ }^{88}$ | $+0.2129 \varphi^{\prime}{ }^{\text {ss }}$ | $+0.1680 \varphi_{\mathrm{N}^{3 s}}$ |
| 2 | -15.441 | 0 | $-0.2565 \varphi_{1}{ }^{\text {sa }}$ | $-0.3068 \varphi_{\varphi_{1}} \mathrm{ss}$ | $-0.3002 \varphi_{\mathrm{N}^{\text {se }}}$ |
| 3 | $-14.975$ | $0.1686 \varphi_{0^{\text {ab }}}$ | $+0.1451 \varphi_{1}{ }^{\text {as }}$ | $+0.3168 \varphi^{19}$ | $+0.3382 \varphi_{\mathrm{N}^{3 s}}$ |
| 4 | $-14.749$ | 0 | $-0.0821 \varphi_{1}{ }^{\text {as }}$ | $-0.3280 \varphi_{\varphi^{198}}$ | $-0.3682 \varphi_{\mathrm{N}^{88}}$ |
| 5 | $-13.742$ | $-0.2585 \varphi_{0}{ }^{\text {ss }}$ | $-0.1940 \varphi_{1}{ }^{38}$ | $+0.24229^{188}$ | $+0.3468 \varphi_{\mathrm{N}^{8 s}}$ |
| 6 | $-12.143$ | $-0.5489 \varphi_{0}{ }^{\text {as }}$ | $-0.2520 \varphi^{9}{ }^{\text {as }}$ | $+0.0768 \varphi^{1 / 88}$ | $+0.1730 \varphi_{\mathrm{N}^{88}}$ |
| 7 | $-11.782$ | 0 | $-0.4171 \varphi_{1}{ }^{38}$ | +0.0958 $\varphi_{1}{ }^{\text {sa }}$ | $+0.2585 \varphi_{\mathrm{N}}{ }^{\text {sa }}$ |
| 8 | -2.318 | 0 | $-0.4029 \varphi_{1}{ }^{\text {sa }}$ | $-0.1707 \varphi_{1}{ }^{\text {sa }}$ | $+0.2419 \varphi_{\mathrm{N}^{38}}$ |
| 9 | $-1.637$ | $-0.5419 \varphi_{0}{ }^{88}$ | $+0.2314 \varphi_{1}{ }^{\text {s }}$ | $+0.1420 \varphi_{1^{188}}$ | $+0.1717 \varphi_{\mathrm{N}}{ }^{\text {sa }}$ |
| 10 | $-0.345$ | $0.2673 \varphi^{\text {a }}{ }^{\text {a }}$ | $-0.1802 \varphi_{1}{ }^{\text {ss }}$ | $-0.3058 \varphi^{188}$ | $+0.2971 \varphi_{\mathrm{N}^{88}}$ |
| 11 | 0.825 | 0 | $-0.1011 \varphi_{1}{ }^{\text {sa }}$ | $+0.3830 \varphi_{1_{1}{ }^{\text {sa }}}$ | $-0.3051 \varphi_{\mathrm{N}}{ }^{\text {es }}$ |
| 12 | 1.155 | $-0.2072 \varphi_{0} 8 \mathrm{ss}$ | $+0.1744 \varphi_{1}{ }^{89}$ | $-0.3551 \varphi_{1^{188}}$ | $+0.2684 \varphi_{\mathrm{N}}{ }^{\text {sa }}$ |
| 13 | 1.721 | 0 | +0.2844 $\varphi_{1}{ }^{88}$ | $-0.3365 \varphi_{1_{1}}{ }^{\text {s3 }}$ | $+0.2363 \varphi_{\mathrm{N}}{ }^{\text {as }}$ |
| 14 | 2.993 | $-0.3143 \varphi_{0}{ }^{\text {a }}$ | +0.3646 $\varphi_{1}{ }^{\text {as }}$ | $-0.2241 \varphi_{1_{1}}{ }^{\text {as }}$ | $+0.1319 \varphi_{\mathrm{N}^{\mathrm{am}}}$ |

${ }^{a} \varphi_{0}{ }^{s \mathrm{~s}}=\chi_{01}+\chi_{02}, \varphi_{0}^{a s}=\chi_{01}-\chi_{02}: \varphi_{1}^{s B}=\chi_{1}+\chi_{2}+\chi_{3}+\chi_{4} ; \varphi_{1}^{s a}=\chi_{1}-\chi_{2}+\chi_{3}-\chi_{4} ; \varphi_{1}^{a s}=\chi_{1}+\chi_{2}-\chi_{3}-\chi_{4} ; \varphi_{1}^{a s}=\chi_{1}-$ $\chi_{2}-\chi_{3}+\chi_{4} ; \quad \varphi_{1}^{\prime}$ consists of $\chi_{1}{ }^{\prime}, \chi_{2}{ }^{\prime}, \chi_{3}{ }^{\prime}$, and $\chi_{4}{ }^{\prime}$ instead of $\chi_{1}, \chi_{2}, \chi_{3}$, and $\chi_{4}$, respectively, in the corresponding $\varphi_{1}$ orbital. $\varphi_{N}$ consists of $\chi_{N_{1}}, \chi_{N_{2}}, \chi_{N_{3}}$, and $\chi_{N_{4}}$ instead of $\chi_{1}, \chi_{2}, \chi_{3}$, and $\chi_{4}$, respectively, in the corresponding $\varphi_{1}$ orbital. ${ }^{b}$ The atoms of TCNB are marked as follows


The locally (in TCNB) excited configurations are given by ${ }^{25}$

$$
\begin{align*}
& \Phi_{\mathrm{L} 1}=a \Phi\left(\phi_{6}{ }^{-1} \phi_{8}\right)+b \Phi\left(\phi_{7}{ }^{-1} \phi_{9}\right) \\
& \quad a=0.7878, b=0.6159 \\
& \Phi_{\mathrm{L} 2}=c \Phi\left(\phi_{6}{ }^{-1} \phi_{9}\right)+d \Phi\left(\phi_{7}{ }^{-1} \phi_{8}\right)  \tag{6}\\
& \quad c=-0.5032, d=0.8642
\end{align*}
$$

where
$\Phi\left(\phi_{i}^{-1} \phi_{k}\right)=1 / 2\left\{\left|\theta_{1} \ldots \phi_{i} \bar{\phi}_{k} \ldots \bar{\phi}_{7}\right|+\left|\theta_{1} \ldots \phi_{k} \bar{\phi}_{i} \ldots \bar{\phi}_{7}\right|\right\}$
The wave functions $\Phi_{L_{1}}$ and $\Phi_{L_{2}}$ can be represented by a linear combination of configurations caused by oneelectron excitation, the coefficients $a, b, c$, and $d$ being determined by the configuration interaction calculation for TCNB itself. The four CT configurations are as follows.

$$
\begin{align*}
& \Phi_{\mathrm{CT}_{1}}=\Phi\left(\theta_{2}^{-1} \phi_{8}\right) \\
& \Phi_{\mathrm{CT}_{2}}=\Phi\left(\theta_{2}^{-1} \phi_{9}\right)  \tag{7}\\
& \Phi_{\mathrm{CT}_{3}}=\Phi\left(\theta_{3}^{-1} \phi_{8}\right) \\
& \Phi_{\mathrm{CT}_{4}}=\Phi\left(\theta_{3}^{-1} \phi_{9}\right)
\end{align*}
$$

In order to consider the configuration interaction among the above-mentioned seven electron configurations, we must evaluate the energies of these configurations and the off-diagonal elements of the total electronic Hamiltonian for the system.
The energies of the excited configurations are evaluated by taking the energy of the ground configurations as the standard. The energies for the $\Phi_{\mathrm{LI}}$ and
(25) Locally excited configuration with higher energies are not considered in the present calculation. In particular it may be noticed that the locally excited configurations of the donor corresponding to the $A_{1 g}$ $\rightarrow B_{2 u}, A_{1 g} \rightarrow B_{l_{u}}$, and $A_{g} \rightarrow E_{1_{u}}$ transitions of benzene are disregarded. This is because the locally excited configurations of donors are higher than the lowest locally excited configuration of the acceptor (TCNB), and because the interaction of the CT configurations with locally excited confgurations of the donor is small in the stable form of the complex, as is revealed from the overlap integral values given in Table VII.
$\Phi_{\mathrm{L} 2}$ configurations are taken from the transition energies observed with TCNB.

$$
\begin{gather*}
E_{\mathrm{G}}=0 \mathrm{ev} \\
E_{\mathrm{L} 1}=3.93 \mathrm{ev}  \tag{8}\\
E_{\mathrm{L} 2}=4.70 \mathrm{ev}
\end{gather*}
$$

The energies of the CT configurations can be estimated as

$$
\begin{align*}
& E_{\mathrm{CT}_{1}}=I-A-C_{28} \\
& E_{\mathrm{CT}_{2}}=I-A^{\prime}-C_{29} \simeq I-A-0.68-C_{29} \\
& E_{\mathrm{CT}_{3}}=I-A-C_{38}  \tag{9}\\
& E_{\mathrm{CT}_{4}}=I-A^{\prime}-C_{39} \simeq I-A-0.68-C_{39}
\end{align*}
$$

where $I$ is the ionization potential of the donor, $A$ and $A^{\prime}$ are the electron affinities for the first and the second lowest vacant orbitals, the difference of which is assumed to be the same as the difference between the SCF orbital energies, $\epsilon_{8}$ and $\epsilon_{9}$. The electron affinity $A$ is taken as one of the parameters. The Coulomb integrals

$$
C_{i j}=\int \theta^{*}{ }_{i}(1) \theta_{i}(1) \frac{e^{2}}{r_{12}} \phi_{j}^{*}(2) \phi_{j}(2) \mathrm{d} \tau_{1} \mathrm{~d} \tau_{2}
$$

are evaluated with the point-charge approximation

$$
\begin{align*}
C_{i j} \cong & \sum_{p} \sum_{q} c_{i p}{ }^{2} c_{j q}{ }^{2} \int \chi_{p}{ }^{*}(1) \chi_{p}(1) \times \\
& \frac{e^{2}}{r_{12}} x_{q}^{*}(2) \chi_{q}(2) \mathrm{d} \tau_{1} \mathrm{~d} \tau_{2} \cong \sum_{p} \sum_{q} c_{i p}{ }^{2} c_{j q}{ }^{2} \frac{e^{2}}{R_{p q}} \tag{10}
\end{align*}
$$

where $x_{p}$ and $x_{q}$ are the atomic orbitals on $p$ and $q$ atoms of donor and acceptor, respectively, and $R_{p q}$ is the distance between them. In eq 9 we have neglected the electrostatic interaction between the donor and the acceptor in the ground configuration.
The off-diagonal terms of the CT configurations with the ground and excited configurations are given in


Figure 9. One of the models assumed for the TCNB-benzene complex.

Table VI. They are given in terms of the following integral

$$
\left(H^{c}\right)_{i j} \equiv \int \theta_{i}(1) H^{\mathrm{c}}(1) \phi_{j}(1) \mathrm{d} \tau
$$

Table VI. The Matrix Elements between the CT Configurations and the Ground and the Locally Excited Configurations

|  | $\Phi_{\mathrm{G}}$ | $\Phi_{\mathrm{L} 1}$ | $\Phi_{\mathrm{L} 2}$ |
| :---: | :---: | :---: | :---: |
| $\Phi\left(\theta_{2}{ }^{-1} \phi_{8}\right)$ | $\sqrt{2}\left(H^{\mathrm{c}}\right)_{28}$ | $a\left(H^{\circ}\right)_{28}$ | $d\left(H^{\circ}\right)_{27}$ |
| $\Phi\left(\theta_{2}{ }^{-1} \phi_{9}\right)$ | $\sqrt{2}\left(H^{\circ}\right)_{28}$ | $b\left(H^{\circ}\right)_{27}$ | $c\left(H^{\circ}\right)_{28}$ |
| $\Phi\left(\theta_{3}{ }^{-1} \phi_{8}\right)$ | $\sqrt{2}\left(H^{\circ}\right)_{38}$ | $a\left(H^{\circ}\right)_{38}$ | $d\left(H^{\circ}\right)_{37}$ |
| $\Phi\left(\theta_{3}{ }^{-1} \phi_{9}\right)$ | $\sqrt{2}\left(H^{\circ}\right)_{39}$ | $b\left(H^{\circ}\right)_{37}$ | $c\left(H^{\circ}\right)_{38}$ |

Here $H^{c}$ is the core attraction term of the total electronic Hamiltonian of the system. The integral is assumed to be proportional to the overlap integral

$$
\begin{equation*}
\left(\theta_{i}\left|H^{\mathrm{c}}\right| \phi_{j}\right)=-K S_{i j} \tag{11}
\end{equation*}
$$

where $S_{i j}=\int \theta_{i}(1) \phi_{j}(1) \mathrm{d} \tau$
The constant $K$ is taken as one of the experimental parameters. The matrix elements between the CT configurations are shown to be very small by an evaluation with the point charge approximation.

For calculating the overlap integrals, ${ }^{26}$ two types of carbon atomic orbital functions were employed: one is the minimal Slater-type orbital with the orbital exponent $1.568^{27}$ and the other, Roothaan's expanded Hartree-Fock atomic orbitals. ${ }^{27}$ The values evaluated by the latter atomic orbitals are two or three times larger than those by the former atomic orbitals. ${ }^{28}$

The calculation was made on the basis of the geometrical models shown in Figure 9, where the center of one benzene ring is taken as an origin of the coordinate system and $R_{x}, R_{y}$, and $R_{z}$ are the coordinates of the center of the other benzene ring. For the sake of simplicity the separation between the two benzene rings is fixed as $R_{z}=3.4 \AA$, and only the displacement along the $x$ and $y$ axes is taken into account.

Model 1; $R_{x}=0$ and $R_{v}=0$. In this case, the overlap integrals $S_{28}, S_{29}, S_{38}$, and $S_{39}$, which are related to the stabilization of the ground state, become zero from the symmetry of the molecular orbitals. ${ }^{29}$ There-

[^4]

Figure 10. The values of the overlap integrals in $R_{y}=0, R_{z}=$ 3.4, and $R_{x}=0-2.0 \AA$.
fore the lower CT configurations do not contribute to the stabilization of the ground state. This seems to mean that the displaced conformation, model 2, may be more favorable for the formation of the stable complex than model 1 .

In model 1 the higher CT configurations interact with the ground configuration through the overlap integrals $S_{19}, S_{2,10}$, and $S_{8,11}$. These overlap integrals are given in Table VII.

Model 2; $R_{x} \neq 0$ or $R_{\nu} \neq 0$. Since the overlap integrals between the occupied orbitals of the donor and the vacant orbitals of the acceptor are an important factor in considering the stabilization of the complex, they were calculated for some typical structures, with the results given in Table VII and Figure 10. The inspection of the results shows that these overlap integrals have the maximum values in the structure where a center of one benzene ring is located above a $\mathrm{C}-\mathrm{C}$ bond of the other ring.

The stabilization energy of the ground state and the transition energies for excited states are tentatively evaluated and compared with the experimental results. The secular equations have been solved for two cases: (a) $R_{x}=1.2035, R_{y}=0.695$, and $R_{z}=3.4 \AA$; (b) $R_{x}=1.2035, R_{y}=0$, and $R_{z}=3.4 \AA$.

The Coulomb terms (ev) in the CT diagonal energies are obtained as follows.
$C_{28}=$ (a) 3.294 (b) $3.326 \quad C_{29}=$ (a) 3.406 (b) 3.439
$C_{33}=$ (a) 3.316 (b) $3.353 \quad C_{39}=$ (a) 3.411 (b) 3.445
We must determine two parameters, $I-A$ and $K$. $I-A$ has been selected in the range from 8.65 to 6.25 ev. When the Hartree-Fock atomic orbitals are used, the value of $K$ is settled to 10 , the magnitude of which might be reasonable from the comparison with the case of benzene, where $\beta_{\mathrm{C}-\mathrm{c}} / S_{\mathrm{C}-\mathrm{C}}=-9.6$. However, by the single Slater-type orbital for calculation between molecules, a much larger $K$ value, 20-30, is used since the overlap integral is underestimated by this procedure.

Table VII. The Overlap Integrals between Benzene and TCNB Orbitals

| $\begin{gathered} R_{x} \\ \mathbf{A} \end{gathered}$ | ${ }_{\AA}^{R_{i}}$ | $R_{\mathrm{A}},$ | $S_{28}$ | $S_{29}$ | $S_{38}$ | $S_{39}$ | $S_{2,10}$ | $S_{3.11}$ | $S_{28}$ | $S_{37}$ | $S_{58}{ }^{\text {a }}$ | $S_{49}{ }^{\text {a }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0 | 0 | 3.4 | 0 | 0 | 0 | 0 | -0.0043 | 0.0133 | -0.0773 | -0.0570 | -0.0234 | -0.0238 |
| 0.40 | 0 | 3.4 | 0.0096 | 0 | 0 | -0.0111 | -0.0052 | 0.0135 | -0.0745 | -0.0509 | -0.0214 | -0.0209 |
| 0.80 | 0 | 3.4 | 0.0177 | 0 | 0 | -0.0185 | -0.0076 | 0.0136 | 0.0664 | -0.0342 | -0.0160 | -0.0135 |
| 1.20 | 0 | 3.4 | 0.0229 | 0 | 0 | -0.0199 | -0.0105 | 0.0125 | -0.0544 | -0.0116 | -0.0082 | -0.0050 |
| 1.60 | 0 | 3.4 | 0.0245 | 0 | 0 | -0.0155 | -0.0126 | 0.0092 | -0.0400 | -0.0113 | 0.0002 | 0.0015 |
| 0 | 0.46 | 3.4 | , | -0.0091 | -0.0110 | 0 | -0.0032 | 0.0133 | -0.0693 | -0.0543 | -0.0200 | -0.0215 |
| 0 | 0.93 | 3.4 | 0 | -0.0144 | -0.0195 | 0 | -0.0006 | 0.0133 | -0.0480 | -0.0469 | -0.0107 | -0.0158 |
| 0 | 1.39 | 3.4 | , | -0.0140 | -0.0239 | 0 | +0.0021 | 0.0130 | -0.0204 | -0.0365 | -0.0006 | -0.0095 |
| 1.20 | 0.695 | 3.4 | 0.0167 | -0.0132 | -0.0051 | -0.0198 |  |  | -0.0417 | -0.0089 |  |  |
| 1.20 | 0.695 | 2.8 | 0.0356 | -0.0275 | -0.0106 | -0.0421 |  |  | $-0.0776$ | $-0.0159$ |  |  |

${ }^{a}$ Overlap integrals concerning the interaction between CT configurations and lower locally excited configurations of benzene.

Table VIII. Calculated Wave Functions and Energies for the Mesitylene-TCNB Complex ${ }^{a}$

|  | $W_{i}, \mathrm{ev}$ |  | $\Phi_{G}{ }^{\text {b }}$ | $\Phi_{L 1}{ }^{\text {c }}$ | $\Phi_{\mathrm{L} 2}{ }^{\text {c }}$ | $\Phi_{\mathrm{CT}_{1}{ }^{\text {d }}}$ | $\Psi_{\mathrm{CT}_{2}{ }^{\text {a }}}$ | ¢ $_{\mathrm{CT}_{3}{ }^{\text {d }}}$ | $\Phi_{\mathrm{CT}_{4}{ }^{\text {a }}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $W_{0}$ | -0.05 | $\Psi$ | 0.9941 | 0.0026 | 0.0052 | 0.0665 | -0.0449 | -0.0201 | -0.0699 |
| $W_{1}$ | 3.28 | $\Psi_{1}$ | -0.0358 | 0.5013 | 0.0837 | 0.7244 | 0.0568 | 0.4595 | -0.0353 |
| $W_{2}$ | 3.49 | $\Psi_{2}$ | 0.0491 | -0.0510 | 0.0072 | -0.5037 | -0.0596 | 0.8588 | 0.0087 |
| $W_{3}$ | 3.96 | $\Psi_{3}$ | 0.0710 | 0.4039 | $-0.3131$ | -0.2619 | 0.7590 | -0.0813 | 0.2872 |
| $W_{4}$ | 4.07 | $\Psi_{4}$ | 0.0451 | 0.0892 | -0.0304 | -0.0324 | -0.4212 | -0.0543 | 0.8987 |
| $W_{5}$ | 4.21 | $\Psi_{5}$ | -0.0199 | 0.7582 | 0.1009 | $-0.3681$ | $-0.4006$ | -0.1955 | -0.2836 |
| $W_{B}$ | 4.80 | $\Psi_{8}$ | 0.0245 | 0.0117 | 0.9401 | -0.1098 | 0.2779 | -0.0552 | 0.1523 |

${ }^{a}$ The calculation was made for the following geometrical parameters and the $I-A$ value: $R_{x}=1.20, R_{y}=0.695, R_{z}=3.4 \AA$, and $I-A$ $=6.80 \mathrm{ev},{ }^{b} \Phi_{\mathrm{G}}$ is the wave function of the ground configuration shown in eq $5 .{ }^{c} \Phi_{1}$ and $\Phi_{2}$ are the wave functions of the locally excited configurations of TCNB shown in eq 6. ${ }^{d} \Phi_{\mathrm{CT}_{1}}, \Phi_{\mathrm{CT}_{2}}, \Phi_{\mathrm{CT}_{3}}$, and $\Phi_{\mathrm{CT}_{4}}$ are the wave functions of the CT configurations shown in eq 7 .

Figure 7 shows the transition energies calculated for case a, the Hartree-Fock atomic orbital being used. The evaluated wave functions of the ground and excited states of the TCNB-mesitylene complex as an example are tabulated in Table VIII.

By the use of the resultant wave functions, we calculated the transition probability between the ground and the excited states. Writing the ground and the excited states as

$$
\begin{align*}
& \Psi_{\mathrm{G}}=a \Phi_{\mathrm{G}}+\Sigma b_{i} \Phi_{i}\left(\mathrm{D}^{+} \mathrm{A}^{-}\right)+\Sigma{c_{i} \Phi_{i}\left(\mathrm{DA}^{*}\right)}^{\Psi_{\mathrm{E}}=a^{\prime} \Phi_{\mathrm{G}}+\Sigma b^{\prime}{ }_{i} \Phi_{i}\left(\mathrm{D}^{+} \mathrm{A}^{-}\right)+\Sigma{c_{i}}^{\prime} \Phi_{i}(\mathrm{DA})^{*}} \tag{12}
\end{align*}
$$

the transition moment is approximately represented as

$$
\begin{gather*}
\left(\Psi_{\mathrm{G}}|e r| \Psi_{\mathrm{E}}\right) \simeq b_{i} b_{i}{ }^{\prime}\left(\Phi_{i}\left(\mathrm{D}^{+} \mathrm{A}^{-}\right)|e n| \Phi_{i}\left(\mathrm{D}^{+} \mathrm{A}^{-}\right)\right)+ \\
\Sigma\left(a c_{i}{ }^{\prime}+c_{i} a^{\prime}\right)\left(\Phi_{i}\left(\mathrm{DA}^{*}\right)|e r| \Phi(\mathrm{DA})\right)+\Sigma\left(a b_{i}{ }^{\prime}+\right. \\
\left.b_{i} a^{\prime}\right)\left(\Phi_{i}\left(\mathrm{D}^{+} \mathrm{A}^{-}\right)|e r| \Phi(\mathrm{DA})\right) \tag{13}
\end{gather*}
$$

where the second term of the right-hand side was taken from the experimental values. The first and the third terms are estimated as

$$
\begin{gather*}
\left(\Phi\left(\mathrm{D}^{+} \mathrm{A}^{-}\right)|e r| \Phi\left(\mathrm{D}^{+} \mathrm{A}^{-}\right)\right) \cong e\left(R_{\mathrm{A}}-R_{\mathrm{D}}\right) \\
\left(\Phi\left(\mathrm{D}^{+} A^{-}\right)|e r| \Phi(\mathrm{DA})\right) \simeq \sqrt{2} \sum_{p q} c_{\mathrm{A} p} c_{\mathrm{D} q} \times \\
\left.\int \chi_{p}(1) e r \chi_{q}(1) d \tau \simeq \sqrt{2} \sum_{p q} c_{\mathrm{A} q} c_{\mathrm{D} q}\left(R_{p}+R_{q}\right) S_{p \mathrm{~d}}\right] / 2 \tag{14}
\end{gather*}
$$

The second term in eq 13 appears only for the $x$ and $y$ directions in the molecular plane, while the first and third terms might have components in the $z$ direction as well as in the $x$ and $y$ directions. In Figures 6a and $6 \mathrm{~b},\left(\Psi_{\mathrm{G}}|e z| \Psi_{\mathrm{E}_{\mathrm{J}}}\right)^{2}$ and $\left(\Psi_{\mathrm{G}}|e x| \Psi_{\mathrm{E}_{j}}\right)^{2}+\left(\Psi_{\mathrm{G}}|e y| \Psi_{\mathrm{E}_{j}}\right)^{2}+$ $\left|\left(\Psi_{\mathrm{G}}|e z| \Psi_{\mathrm{E}}\right)\right|^{2}$ are shown for each transition.

In view of the ambiguity in the geometrical structure of the TCNB complexes under consideration, too much importance should not be attached to the present theoretical results, especially to the calculated stabilization energy. However, they might be important in giving a qualitative or semiquantitative explanation to the observed phenomena, in particular, to the appearance of the two CT bands, the decrease in the absorption intensities of the first bands of the TCNB complexes with the decrease in the ionization potentials of electron donors, and the polarization of the first and second CT bands.

## Appendix

By the aid of the electronic computer program of SCF-CI calculation based on the Pariser-Parr-Pople approximation, ${ }^{30,31}$ we calculated the molecular orbitals and the transition energies of TCNB and some other polycyanobenzenes. The results for $\beta_{\mathrm{CN}}=-3.28$ ev and $\beta_{\mathrm{CC}}=-2.39 \mathrm{ev}$ are given in Tables V and IX.

Table IX. Calculated and Observed Transition Energies and Oscillator Strengths

\left.| Transition |
| :---: | :---: | :---: | :---: |
|  |$\right)$

(30) R. Pariser and R. G. Parr J. Chem. Phys., 21, 466 (1953).
(31) J. A. Pople, Proc. Phys. Soc. (London), A68, 81 (1955).


[^0]:    (11) This geometrical orientation may be thought to be supported by the observation that the refractive index of the crystal is greater in the direction perpendicular to the elongated axis than in the direction parallel to it.

[^1]:    (14) Strictly speaking, this assumption is not correct. The equilibrium constant of the benzene-TCNB complex given in Table IV shows that the complex formation amounts to $\sim 80 \%$ when TCNB is dissolved in pure benzene. However, this gives only a small effect to the evaluated $\epsilon_{\text {max }}$ value because the free TCNB molecule has also an absorption in the wavelength region under consideration.

[^2]:    (15) L. E. Orgel and R. S. Mulliken, J. Am. Chem. Soc., 79, 4839 (1957).
    (16) J. N. Murrell, ibid., 81, 4839 (1959).

[^3]:    (23) K. S. Mulliken, J. Chim. Phys., 61, 20 (1963).
    (24) H. Hosoya and S. Nagakura, Bull. Chem. Soc. Japan, 37, 249 (1964).

[^4]:    (26) Calculations of the overlap integrals and the secular equations were performed by the aid of an electronic computer FACOM 202 in our institute. The authors acknowledge the assistance of Mr. H. Takahashi for making the program.
    (27) T. L. Gilbert, private communication.
    (28) Through our calculation we have used common atomic orbitals for constructing the vacant and occupied molecular orbitals, because the shape of the vacant $\pi$ orbital may be less deformed in the case of $\pi$ complexes than in the case of the acceptor orbital of iodine complexes. 16,16
    (29) Even if the rotation of one of the benzene rings around the $Z$ axis is allowed, $S_{28}, S_{29}, S_{38}$, and $S_{39}$ are still equal to zero, because the orbitals $\theta_{2}$ and $\theta_{3}$ of benzene are doubly degenerate.

